

Computational Semiclassics

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$$\psi:\mathbb{R}_t\times\mathbb{R}^d_x\rightarrow\mathbb{C}^M$$

with

$$d \gg 1 \quad \text{and} \quad \|\nabla_{t,x}\psi\| \gg 1$$

Self-adjoint operator \hat{H}

Schrödinger equation: $i\varepsilon\partial_t\psi = \hat{H}\psi$

Unitary propagator $U_t = e^{-i\hat{H}t/\varepsilon}$

$$0 < \varepsilon \ll 1$$

Energy function H smooth, subquadratic

Hamilton equation: $\dot{z}_t = \Omega \nabla H(z_t)$

Symplectic flow $\Phi_t : \mathbb{R}^{2d} \rightarrow \mathbb{R}^{2d}$

Herman–Kluk propagator

cf. Herman/Kluk '81, Kay '06,

Swart/Rousse '09,

L/Sattlegger '16

$$\psi = (2\pi\varepsilon)^{-d}\int\limits_{\mathbb{R}^{2d}} \langle g_z,\psi\rangle g_z\,dz$$

$$g_z(x) = (\pi\varepsilon)^{-d/4}\,\exp\!\left(-\tfrac{1}{2\varepsilon}|x-q|^2 + \tfrac{i}{\varepsilon}p\cdot(x-q)\right)$$

$$\begin{aligned} U_t\psi &= (2\pi\varepsilon)^{-d}\int\limits_{\mathbb{R}^{2d}}\langle g_z,\psi\rangle \,U_tg_z\,dz \\ &\stackrel{!}{\approx} (2\pi\varepsilon)^{-d}\int\limits_{\mathbb{R}^{2d}}\textcolor{red}{u_t}(z)\langle g_z,\psi\rangle\,\textcolor{blue}{e}^{iS_t(z)/\varepsilon}g_{z_t(z)}\,dz \end{aligned}$$

$$\textcolor{red}{u_t}(z)=\sqrt{2^{-d}\det\left(\partial_q q_t(z)+\partial_p p_t(z)+i[\partial_p q_t(z)-\partial_q p_t(z)]\right)}$$

$$I_t = (2\pi\varepsilon)^{-d} \int_{\mathbb{R}^{2d}} u_t(z) \langle g_z, \bullet \rangle e^{iS_t(z)/\varepsilon} g_{z_t(z)} dz$$

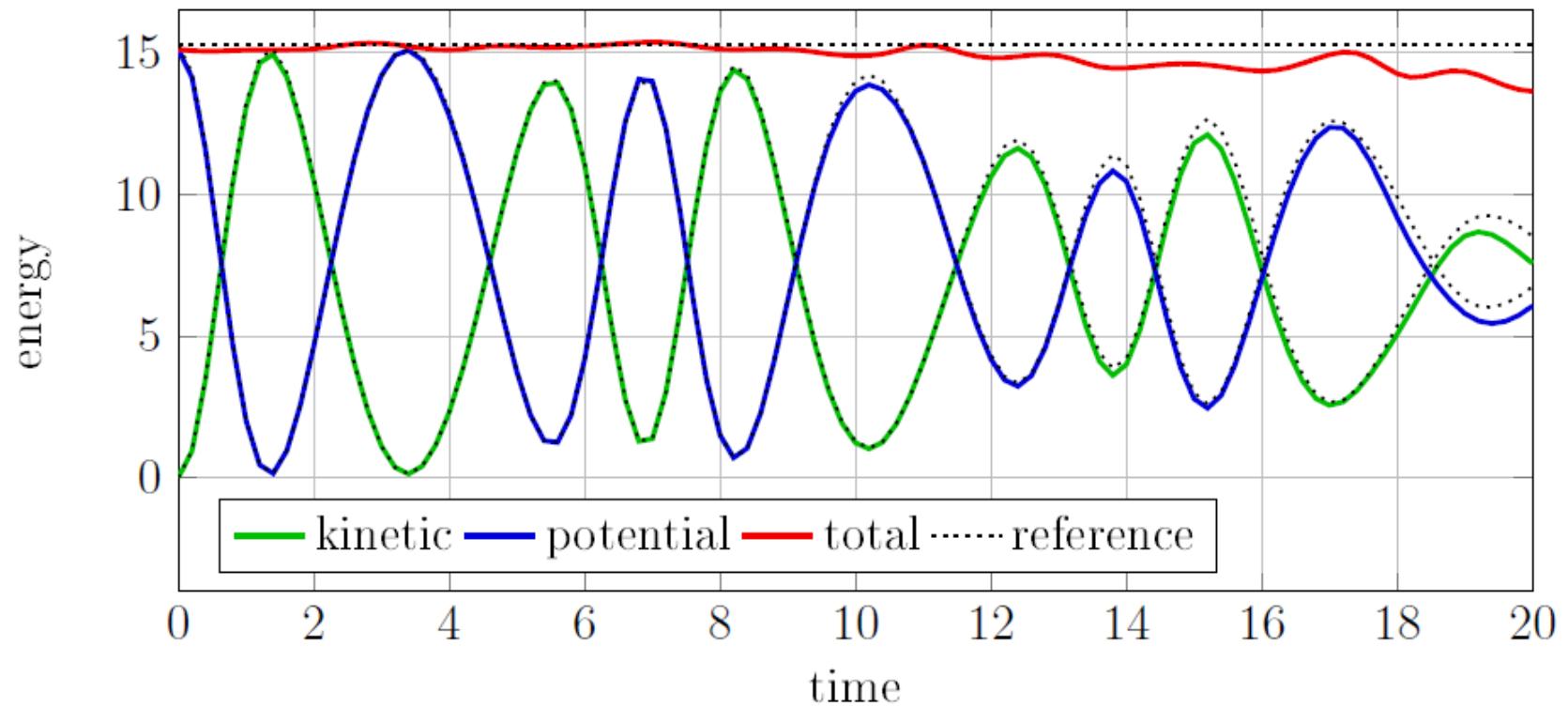
satisfies $\|I_t - U_t\| \leq C_{t,H} \varepsilon$.

If $|\widetilde{\Phi}^\tau - \Phi^\tau| \leq c \tau^\gamma$ is **symplectic**, then $\|I_\tau - \tilde{I}_\tau\| \leq \tilde{C}_{c,H} \tau^\gamma / \varepsilon$.

$$\begin{aligned} I_t\psi &= (2\pi\varepsilon)^{-d}\int\limits_{\mathbb{R}^{2d}} u_t(z)\langle g_z,\psi\rangle \, \mathrm{e}^{iS_t(z)/\varepsilon}g_{z_t(z)}\, dz \\ &\approx \frac{1}{N}\sum_{j=1}^N \widetilde{u}_t(z_j) \, \textcolor{red}{r}_\psi(z_j) \, \mathrm{e}^{i\widetilde{S}_t(z_j)/\varepsilon}g_{\widetilde{z}_t(z_j)} \end{aligned}$$

with $z_1, \dots, z_N \sim \mu_\psi$.

$$d = 6, \varepsilon = 0.01$$



computing time on a desktop computer: 5 resp. 30 minutes

LSC-IVR

cf. Miller '74, Lee & Scully '80, Heller '81,
..., L & Röblitz '10

Egorov's theorem,

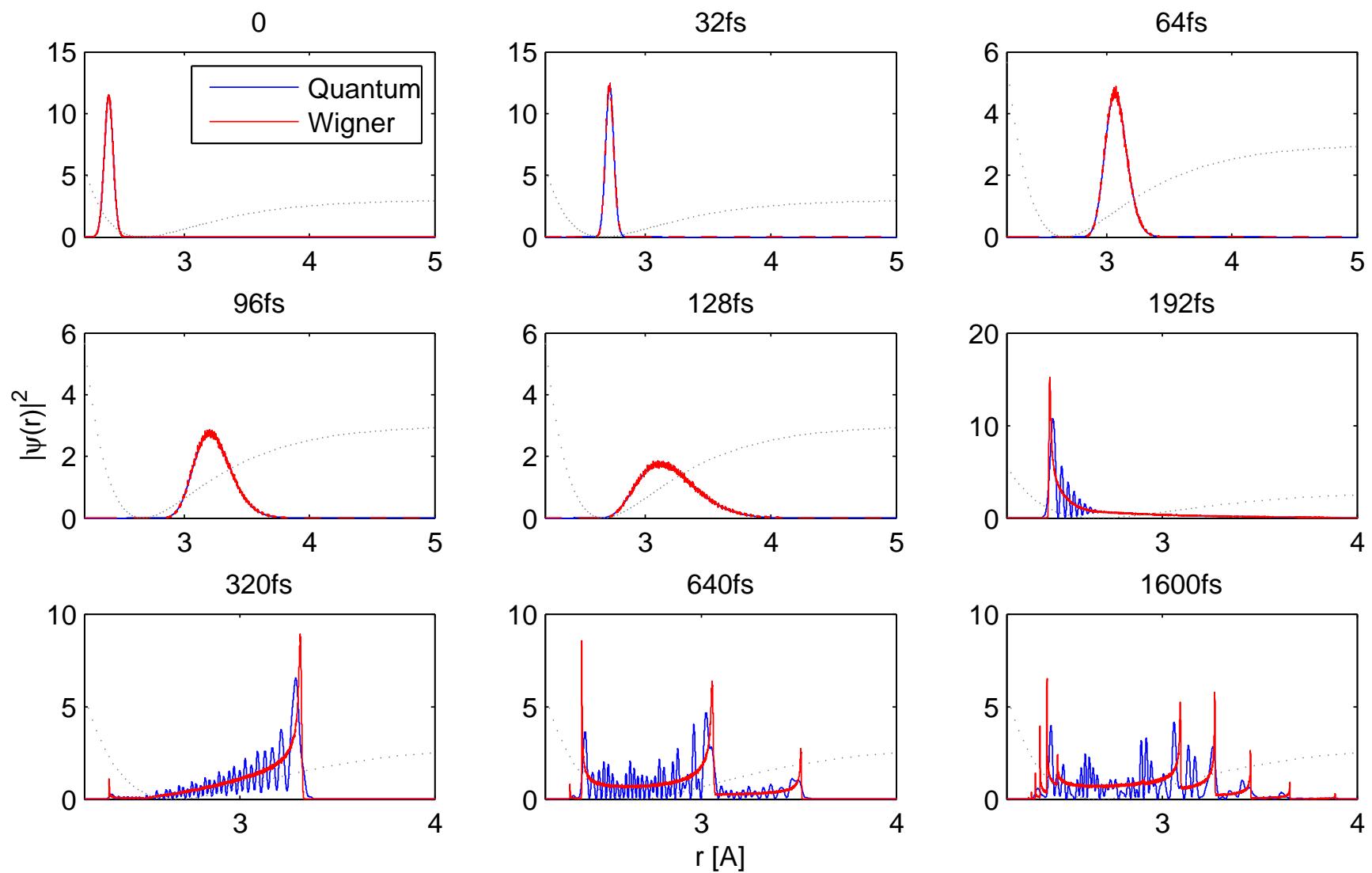
$$\|U_t^* \widehat{A} U_t - \widehat{A \circ \Phi_t}\| \leq C_{t,A,H} \varepsilon^2,$$

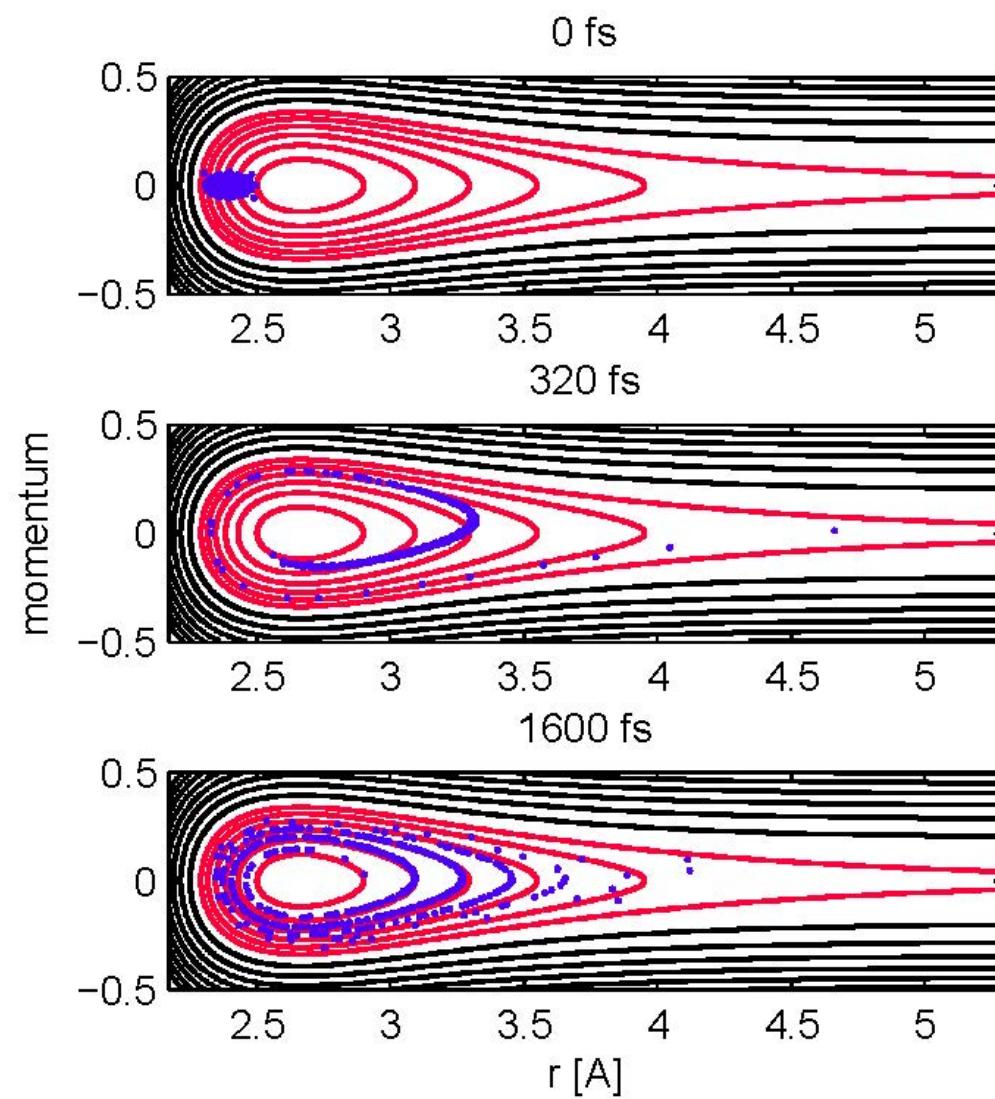
implies

$$\langle U_t \psi, \widehat{A} U_t \psi \rangle \approx \langle \psi, \widehat{A \circ \Phi_t} \psi \rangle.$$

$$\begin{aligned}\left\langle \psi, \widehat{A \circ \Phi_t} \psi \right\rangle &= \int_{\mathbb{R}^{2d}} (A \circ \Phi_t)(z) W_\psi(z) dz \\ &\approx \frac{1}{N} \sum_{j=1}^N (A \circ \widetilde{\Phi}_t)(z_j)\end{aligned}$$

with $z_1, \dots, z_N \sim W_\psi$.



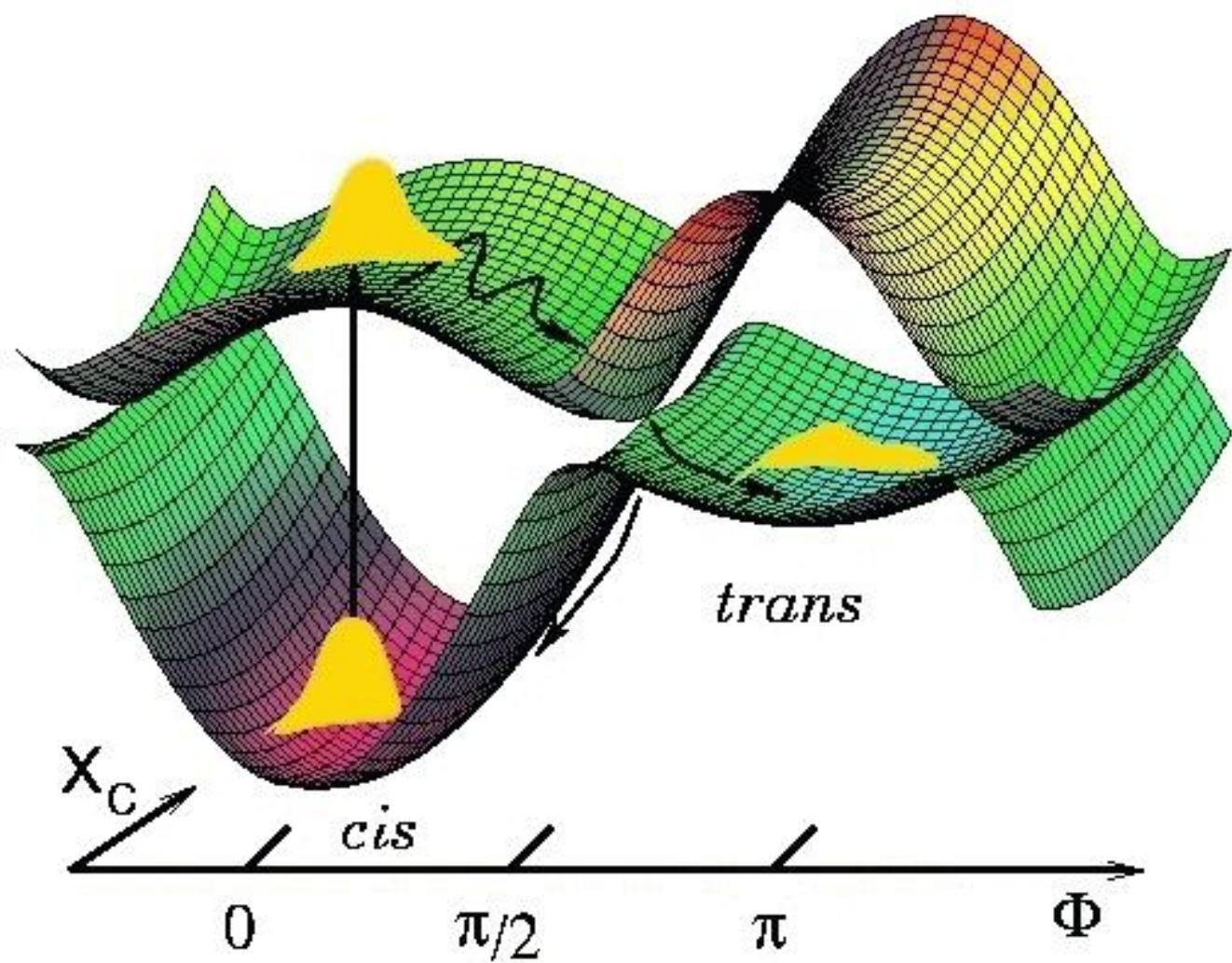


Surface hopping

cf. Bjerre/Nikitin '67, Tully/Preston '71,

Voronin/Marques/Varandas '98,

L/Teufel/Fermanian '03 –'16



Let $H = \lambda^+ \Pi^+ + \lambda^- \Pi^-$.

If $A = \Pi^\pm A \Pi^\pm$ and (\dots) , then

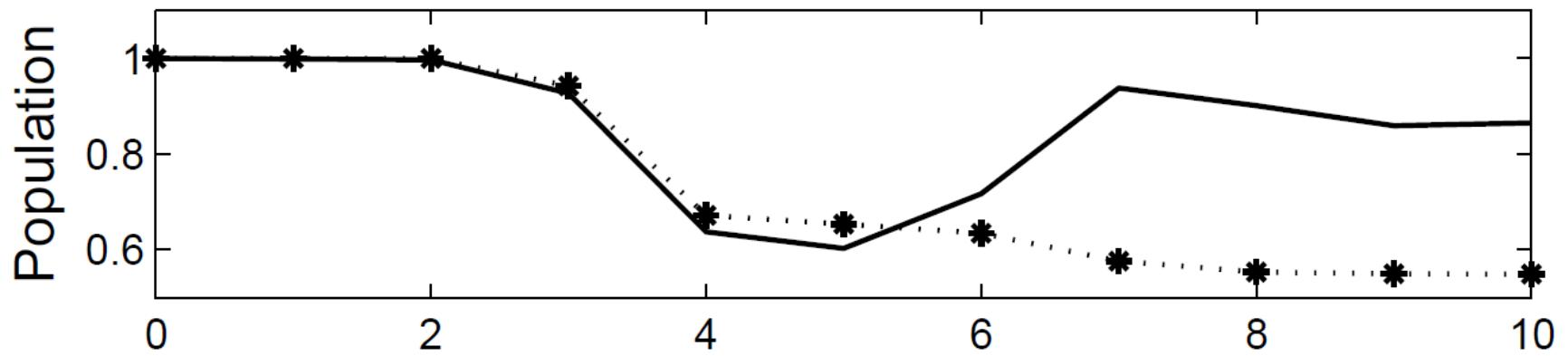
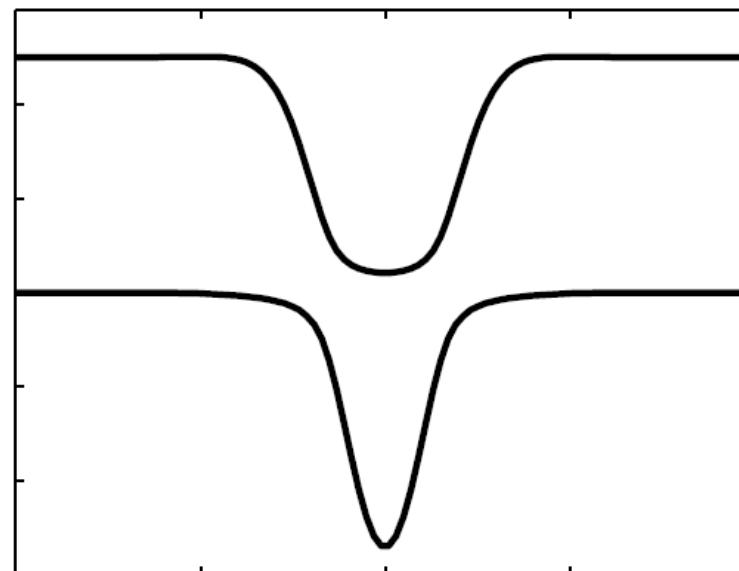
$$\left| \int_0^T \left(\langle U_t \psi, \widehat{A} U_t \psi \rangle - \int_{\mathbb{R}^{2d}} (\mathcal{L}_t A)(z) W_\psi(z) dz \right) dt \right| \leq C_{T,A,H} \varepsilon^{1/8}.$$

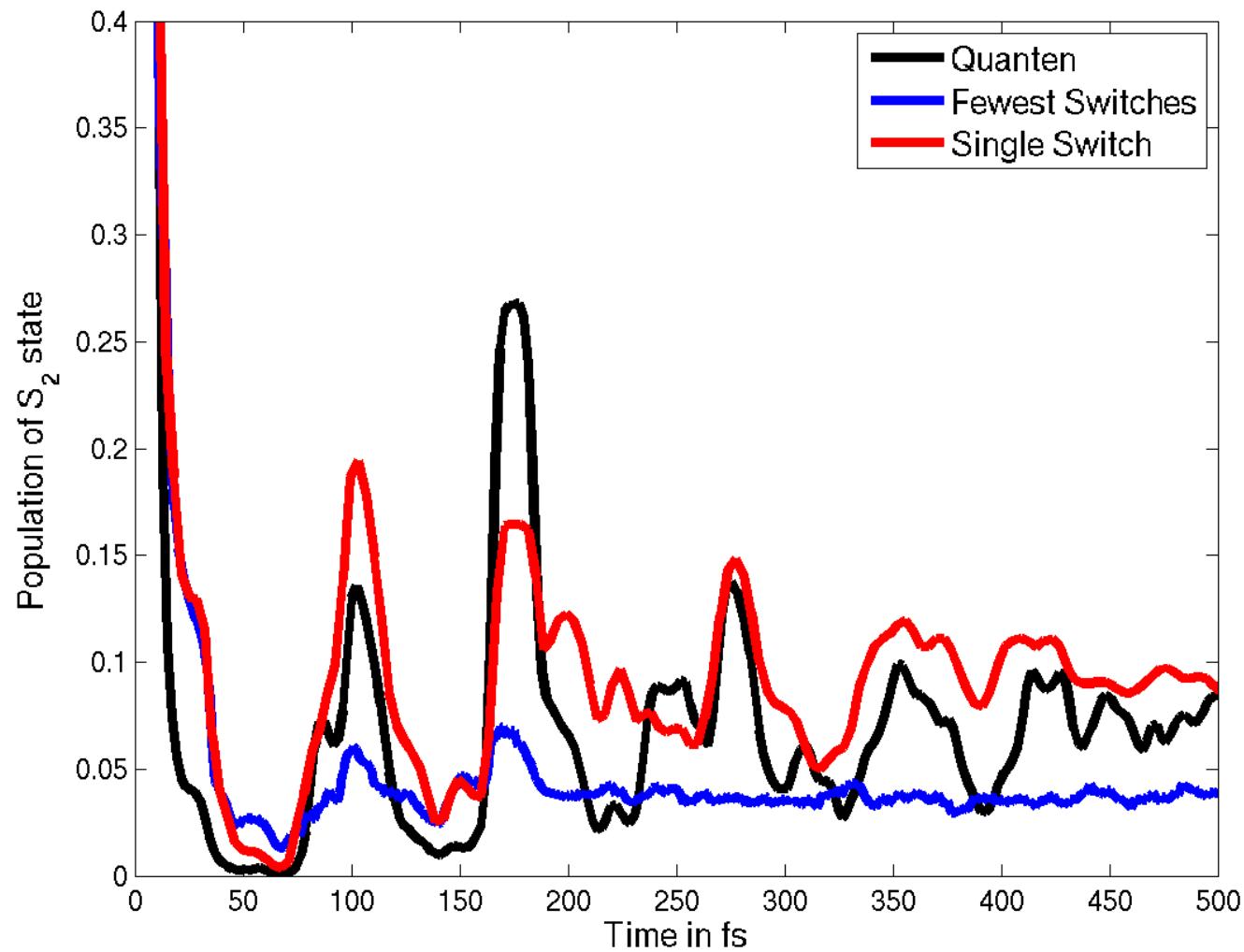
The semigroup \mathcal{L}_t combines

- ⊕ classical transport along λ^+ and λ^- ,
- ⊕ hops between the surfaces, when $t \mapsto (\lambda^+ - \lambda^-)(z_t)$ is minimal.

The hopping probability comes from a Landau–Zener formula.

Dual





Computational semiclassics

(Summary)

1. Herman–Kluk propagator
2. LSC-IVR (Egorov's theorem)
3. Surface hopping